BEACH EROSION BOARD
OFFICE OF THE CHIEF OF ENGINEERS

WAVE FORECASTING RELATIONSHIPS FOR THE GULF OF MEXICO

TECHNICAL MEMORANDUM NO. 84



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DECEMBER 1956

FOR EWORD

The Beach Erosion Board is issuing a series of reports providing wave statistics for selected regions of the coastal waters of the United States. Coverage of the coastal waters of Gulf of Mexico is being provided by a series of five reports, issued as Beach Erosion Board Technical Memoranda 85-89. The need for such data is evident, and it is planned ultimately to supply it by actual wave measurements for sufficiently long periods to establish the wave climate at many locations. However, with present instrumentation and funds, and the long period of record needed to supply adequate statistical data, such a field program for complete shore line coverage does not appear feasible at the present time. The compilation of wave statistics by "hindcast" techniques, though not as exact as recorded data, nevertheless will provide the engineer with better wave data than have heretofore been available.

This report presents the methods utilized in obtaining hindcast statistical wave data for locations in the Gulf of Mexico. It demonstrates a numerical method of forecasting wave generation and propagation over a sloping bottom taking into account both generation by the wind and dissipation by bottom friction. It also shows the averaging techniques used to apply this method to statistical accumulation of hindcast data in the Gulf.

This report was prepared by Charles L. Bretschneider, a hydraulic engineer in the Research Division of the Beach Erosion Board under the supervision of Joseph M. Caldwell, Chief of the Division. At the time the report was prepared, Brigadier General Theron D. Weaver was President of the Board.

Views and conclusions stated in the report are not necessarily those of the Beach Erosion Board.

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WAVE FORECASTING RELATIONS HIPS FOR THE GULF OF MEXICO

by

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INTRODUCTION

This report describes the development and application of a method for computing wind wave data over the Continental Shelf along the United States coast of the Gulf of Mexico. Methods for forecasting deep-water wind waves, and wind waves in shallow water of constant depth, such as shallow lakes and bays, are already available (8, 9, 12,16)*. The problem of the Continental Shelf of the Gulf of Mexico, however, is more complex, sinds here the waves may generate originally in deep water, but then, under the continued influence of wind, are propagated shoreward over the Continental Shelf where bottom friction and type of bottom profile become important. Because of the variety of initial deep-water conditions and bottom profiles, each location requires separate treatment. A set of generalized forecasting curves is required for each location and each direction to bring the waves in over the shallow sloping bottom to the desired depth. Using deep-water forecasting relationships and taking bottom friction into account, a generalized set of dimensionless forecasting relationships is prepared for each of the five locations (Figure 1) in the Gulf of Mexico for which statistical deep-water wave data are compiled (1, 2, 3, 4, 5). The forecasting curves are intended for the most frequent minimum fetch and corresponding wind speed for various deep-water wave height ranges and average bottom conditions of various directions. For the cases of winds parallel to the coast or from land to sea the curves are applicable to all water depths. However, for the case of winds blowing from sea toward land, the forecasting relationships are satisfactory only for depths of about 20 feet or greater, although the technique has been stretched to a depth of 12 feet for cases where winds are not too high. At depths of about 20 feet or less the bottom slope changes too rapidly for the theory to apply, and longer period swell will be breaking in the surf zone, thereby obscuring the wind wave pattern. Examples on the use of the forecasting graphs are given in Appendix A.

It must be emphasized that a certain amount of averaging is required to simplify the problem, and the forecasting curves must be used with discretion for other than marby locations, where refraction becomes important. The forecasting curves result in wave heights to within +10 percent of the average condition, except in the extreme case of shoal areas.

^{*}Numbers in parentheses denote references listed at end of text.

Statistical distributions of wind waves from records (not yet completely analyzed) obtained in the Gulf of Mexico agree fairly closely to those presented by Putz (6) based on wave records for the Pacific Coast of the United States. Locations selected for developing forecasting relationships and for which statistical wave data are computed are shown in Figure 1. Wave statistics for these locations are given in the following Beach Erosion Board publications.

Technical Memorandum No.
85
86
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BASIC EMPIRICAL RELATIONSHIPS FOR WIND-WAVE GENERATION

Empirical relationships for the generation of wind waves may be arrived at by use of the Buckingham Pi-theorem(7), supplemented by wave observations, provided the observations extend over a wide enough range of the variables. In general the variables to consider are as follows:

		Dimension.
Variable	Symbol .	(on the f, 1, t system)
Cross course	**	17.
wind speed	U	1/t
wave height	H	1
wave period	T	t
wave length	L	1
wave speed	С	1/t
fetch length	F	1
wind duration	t	t
acceleration of gravity	g	1/t ²
depth of water	d	1
bottom slope	8d/8x	1/1
bottom friction factor	f	dimensionless
bottom percolation factor	p	1 ² /t
refraction factor	κ_{R}	dimensionless

In setting up the dimensionless relationships one can make use of the Airy wave theory (eq. 1 and 2) to eliminate the variables L and C.

$$L = \frac{gT^2}{2\pi} \tanh \frac{2\pi i}{L} \tag{1}$$

$$C^2 = \frac{gL}{2\pi} \tanh \frac{2\eta \, \underline{1}}{L} \tag{2}$$

S

Thus one obtains equations 3 and 4:

$$H = H (g, U, F, t, d, \delta d/\delta x, f, p, K_{-})$$
(3)

$$T = T (g, U, F, t, d, \delta d/\delta x, f, p, K_)$$
 (4)

By use of the Pi-theorem it can be shown that:

$$gH/U^{2} = \phi \left(\frac{gF}{u^{2}}, \frac{gt}{u}, \frac{gd}{u^{2}}, \delta d/\delta x, f, p, K_{r} \right)^{*}$$
(5)

$$\frac{gT}{U} = \psi \left(\frac{gF}{U^2}, \frac{gt}{U}, \frac{gd}{U^2}, \delta d/\delta x, f, p, K_r \right)$$
 (6)

Equations 5 and 6 can be reduced to certain special cases, some of which are as follows:

a. Deep water, unlimited wind duration:

$$gH_{o}/U^{2} = \phi_{1}\left(\frac{g|_{i}}{U};\right) \tag{5a}$$

$$gT/U = \psi_1 \left(\frac{gF}{U^2} \right) \tag{6a}$$

b. Deep water, unlimited fetch length:

$$gH_0/U^2 = \phi_2\left(\frac{|t|}{U}\right) \tag{5b}$$

$$gT/U = \psi_2 \left(\frac{gt}{U} \right)$$
 (6b)

c. Shallow water, impermeable flat bottom, constant bottom friction factor, and unlimited wind duration:

$$gH/U^2 = \phi_3 (g!/U^2, gd/U^2)$$
 (5c)

$$gT/U = \psi_3 (gF/U^2, gd/U^2)$$
 (6c)

d. Shallow water, impermeable flat bottom, constant bottom friction factor, and unlimited fetch length:

$$gH/U^2 = \phi_A \ (|t/U, gd/U^2)$$
 (5d)

$$gT/U = \Psi_4 \quad (gt/U, gd/U^2)$$
 (6d)

e. Shallow water, impermeable flat bottom, constant bottom friction factor, unlimited fetch and unlimited wind duration:

$$gH/v^2 = \phi_5 (gd/v^2)$$
 (5e)

$$gT/U = \psi_5 (gd/U^2)$$
 (6e)

Expressions 5a, 5b, 6a, and 6b represent the deep-water wave generation relationships (Figure 2) of Bretschneider(8); expressions 5c, 5d, 5e, 6c, 6d, and 6e are strictly for shallow water with a bottom of constant depth, such as shallow lakes and bays.

Equations 5 and 6, applicable to the Continental Shelf of the Gulf of Mexico, are somewhat involved. One can simplify the work by considering the following:

- a. In most cases wind waves in shallow water are limited either by the fetch length or the water depth at the end of the fetch, or a combination of a fetch length and water depth. Thus by introducing the concept of minimum fetch length the parameter gt/U can be eliminated.
- b. If one limits the development to average conditions, the refraction coefficient k can be neglected temporarily, but applied later. This is perhaps more justifiable since it is difficult to determine refraction of short-crested waves of growing period under the continued influence of wind.

Assuming then that the bottom friction factor is constant, one obtains for the generation of wind wares over the Continental Shelf:

gH'
$$/U^2 = \phi_6 \text{ (gF/U}^2, \text{ gd/U}^2, \text{ bottom profile)}$$
 (5f)

$$gT/U^2 = \psi_6 (gF/U^2, gd/U^2, b) ttom profile)$$
 (6f)

where d is the depth of water at station of interest, and F is the minimum fetch from deep water to d/L = 0.5.

The term H', equivalent deep-waler significant wave height*, has been introduced in place of the altual significant wave height H, which can be obtained from:

$$H_{s} = H^{0} \quad K_{s}, \text{ where}$$
 (7)

K is the shoaling factor given by Figure 7, which was prepared from tables by Wiegel(10).

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^{*}The term H° o equivalent deep-water significant wave height, is defined as that height the shallow-water significant wave would have had in deep water if unaffected by refraction or energy losses or gains.

The introduction of H' simplifies numerical computations required to obtain the wave height in shallow water.

Theoretical Considerations. The discussion following is limited to that necessary for the understanding of a numerical procedure for computing wind waves in shallow water. Sverdrup and Munk(11) have presented a detailed theoretical development for generation of wind waves in deep water. The basic differential equation for the steady state condition is:

$$\frac{C}{2} \frac{\partial E}{\partial x} + \frac{E}{2} = \frac{\partial C}{\partial x} = R_t + R_1, \text{ where}$$
 (8)

C = deep-water wave speed $E = 1/8 \rho gH^2$, wave energy

 $\rho = density$

g = acceleration of gravity

H = wave height

R = rate at which energy is being transferred from wind to wave by tangential stress

R_n = rate at which energy is being transferred from wind to wave (or wave to wind) by normal pressure

$$R_{t} + R_{n} = R_{H} + R_{c}, \text{ where}$$
 (9)

R_u = rate at which energy is transformed into wave height, and

R = rate at which energy is transferred into wave speed.

In a report by Bretschneider and Reid(12) a detailed development is presented for the change in wave height (of swell travelling over a shallow bottom) due to bottom friction, percolation and refraction, assuming no wind effect. The basic differential equation for the steady state condition is:

$$nc \frac{\partial (bE)}{\partial x} + Eb \frac{\partial (nc)}{\partial x} = -\left[D_f + D_p\right]b \qquad (10)$$

where

 $B = 1/8 \rho gH^2$, wave energy

C = wave speed

n = fraction of energy travelling with wave speed

b = perpendicular distance between two wave rays

 D_f = rate at which energy is being lost due to bottom friction D_D^f = rate at which energy is being lost due to percolation

Special solutions of equation 10 are given in reference (12).

If the waves are small and consist of a uniform train being generated in shallow water one may write the following differential equation for the steady state condition, neglecting refraction and percolation:

$$\frac{d(ncE)}{dx} = R_H + R_C - D_f \tag{11}$$

Two general cases for equation 11 exist: (a) for D = 0, which results in the deep-water wave forecasting relationships, Figure 2: and (b) for R, + R = 0 which results in the relationship for wave energy loss (wave height reduction) due to bottom friction, Figure 3. A general solution of the above equation is difficult, because one needs to know how RH and Rc are affected by Df. For this report it is assumed that RH and Rc are related to C/U in some manner (the exact manner is not important) such as that given by Sverdrup and Munk(11); and wind speed, U, and minisum fetch length, Pmin, are given by Figure 2. This is not strictly correct but adjustments are made in the selection of an apparent bottom friction factor which permits a computational procedure and gives agreement between observed and computed wave heights under known conlitions. Furthermore, because of the nature of the problem, the closer to steady state conditions in shallow water, the less dependent s the wave period on energy exchange from wind to wave to dissipation by bottom friction. Using the above relationships and the relationships for wave energy loss by bottom friction, Figure 3, it is possible to solve equation 11 by numerical means or successive approximations. An actual mathematical solution of equation 11 is not necessary since various coefficients must be determined from empirical dath, and since in general numerical means are required for a bottom of variable depth and slope. The proper procedure then is to set up a computational procedure and use observations where available, such that a comparison can be made between observations and computations. If satisfactory agreement is reached where wave data are available, then the method can be extended to locations where wave data are absent. This has been done for the special case of a flat bottom using a bottom friction factor f = .01. The resulting relationships agree quite well with wave observations from Lake Okeechobee, Florida, and the nearly flat areas of the Gulf of Mexico, as shown in Figure 4 reproduced in part from Bretschneider(13), The numerical procedures (8) used in arriving at these relationships are similar to those used on the Continental Shelf, discussed below.

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*Minimum fetch and minimum duration o wind are definitions pertaining to steady state conditions. Minimum uration is that duration of wind (equal to or less than actual duration) for any particular wind speed and fetch length required for steady state generation. Minimum fetch length is that fetch length (equal to or less than actual fetch) for any particular wind speed and wind duration equired for steady state generation.

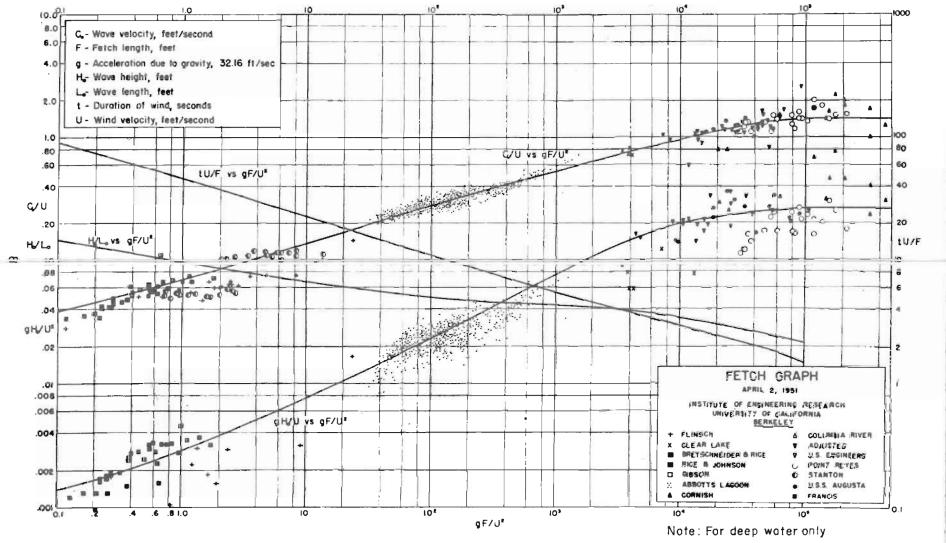


FIGURE 2

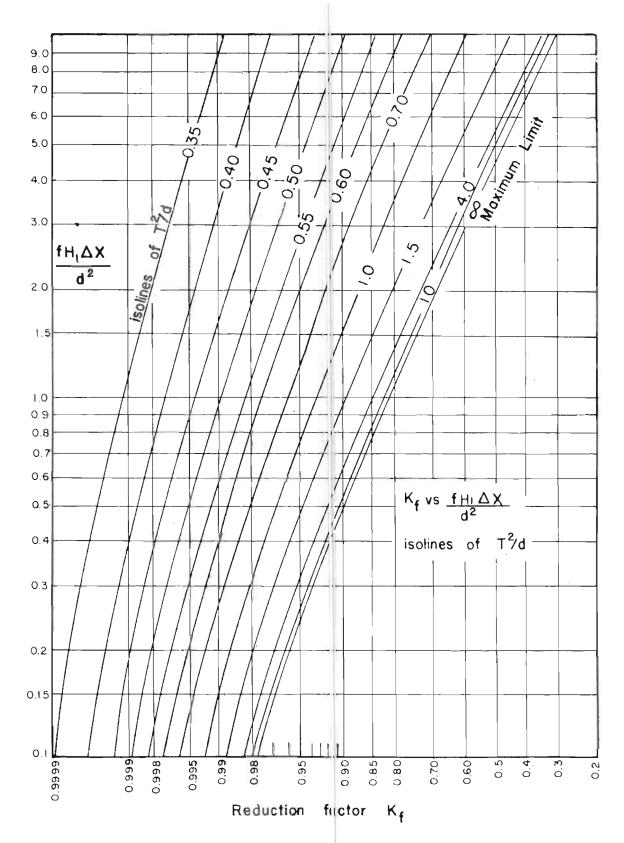


FIGURE 3. RELATIONSHIP FOR FRICTION LOSS OVER

A BOTTOM OF CONSTANT DEPTH

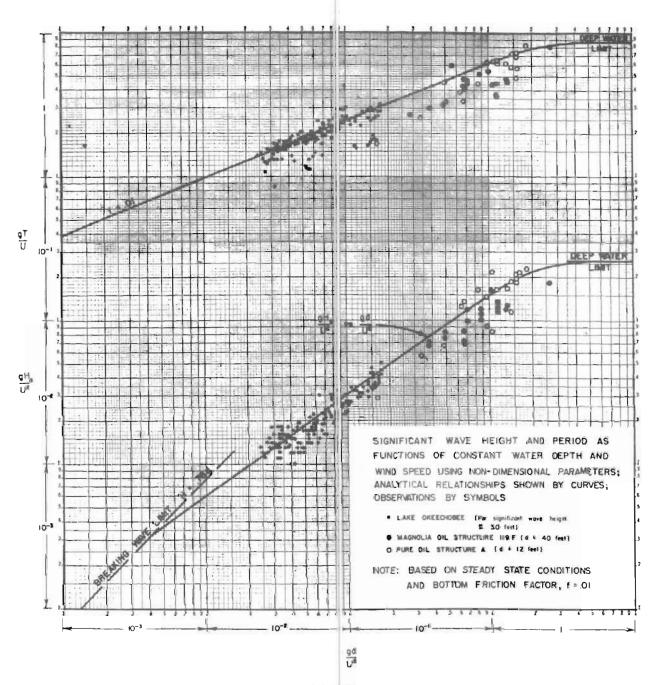


FIGURE 4

Numerical Procedure. A procedure for computing wind waves over a bottom of constant depth and also over a bottom of constant slope has previously been described(9). In general the procedure for a bottom of constant depth can be applied to wind a blowing parallel to the coast and the procedure for a bottom of constant slope can be applied to winds blowing perpendicular to the coast. Through continued use of the above techniques, it was found that the procedure for bottom of constant depth could be used satisfactorily over a bottom of constant slope, provided the bottom slope was segmented average constant depth. Computations ship are much less involved than for sloping bottom. For this reason the bottom profiles are segmented into increments of constant depth.

Figures 2 and 3 are those required to predict the changing wave characteristics as the wave moves over the shallow slope. Figure 2 is used to determine curves of $H_{\rm O}$ and $\Gamma_{\rm S}$ versus fetch length for various wind speeds, and in effect represents wind-wave energy generation. Figure 3 is used to compute a reduction factor $K_{\rm f}$ to be applied to $H_{\rm O}$ to account for the dissipation of vave energy due to bottom friction.

In regard to wave generation by conshore winds there are two conditions to consider. First, the initial deep-water waves generated may be propagated shoreward as swell to nder the continued influence of the generating winds; and second, regeneration of wind waves is constantly taking place all along the fet chover the Continental Shelf. Swell will feel bottom far from shore and commence losing energy at an early stage, whereas wind waves with a shorter period continue to grow and do not feel bottom until they are nearing the coast. In the breaker zore one would observe both swell and wind waves, and the present method of computation cannot be applied here. The steps are as follows: *

- a. A wind speed is selected, and a graph of H_O and T versus fetch distance is computed from Figure 2. Figure 5 is an example for a 26-knot wind and unlimited wind duration.
- b. The minimum fetch length, F_{min} , is selected and H_0 and T is determined at $F = F_{min}$. The deep-water wave length is given by $L_0 = 5.12 \, T^2$, and the waves will begin to feel bottom at a depth $d = L_0/2$. This is the initial point from which to begin computations shoreward.
- c. The bottom profile along the fetch toward the location of interest is determined. The traverse is segmented into at least ten equal increments, ΔF , of about 5 to 10 miles in length, depending on the bottom slope and width of Continental Shelf. Figure 6 is a typical example.
- d. An average depth, dave, is determined over each increment (Figure 6).

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^{*} See Bretschneider (20) for an alternative method for computing wind waves for high wind speeds and short fetches.

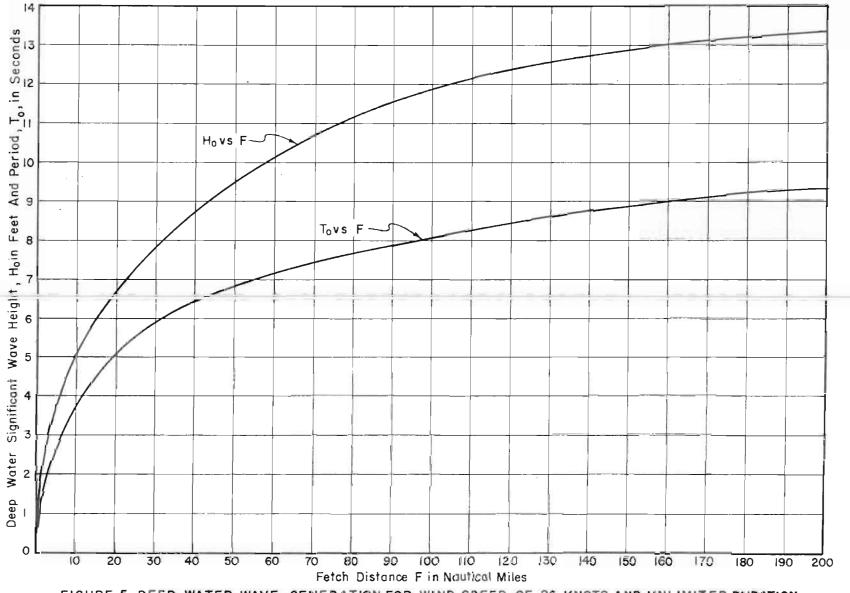
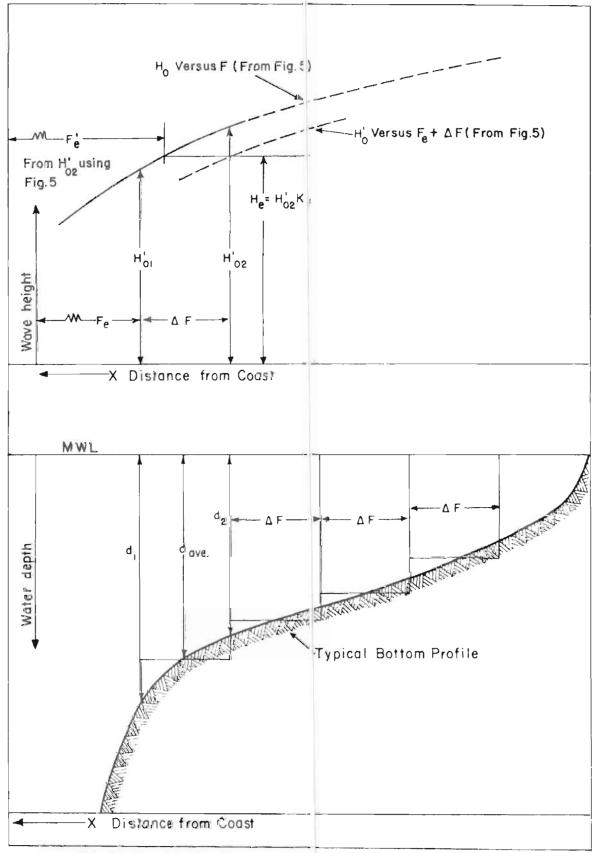


FIGURE 5. DEEP WATER WAVE GENERATION FOR WIND SPEED OF 26 KNOTS AND UNLIMITED DURATION



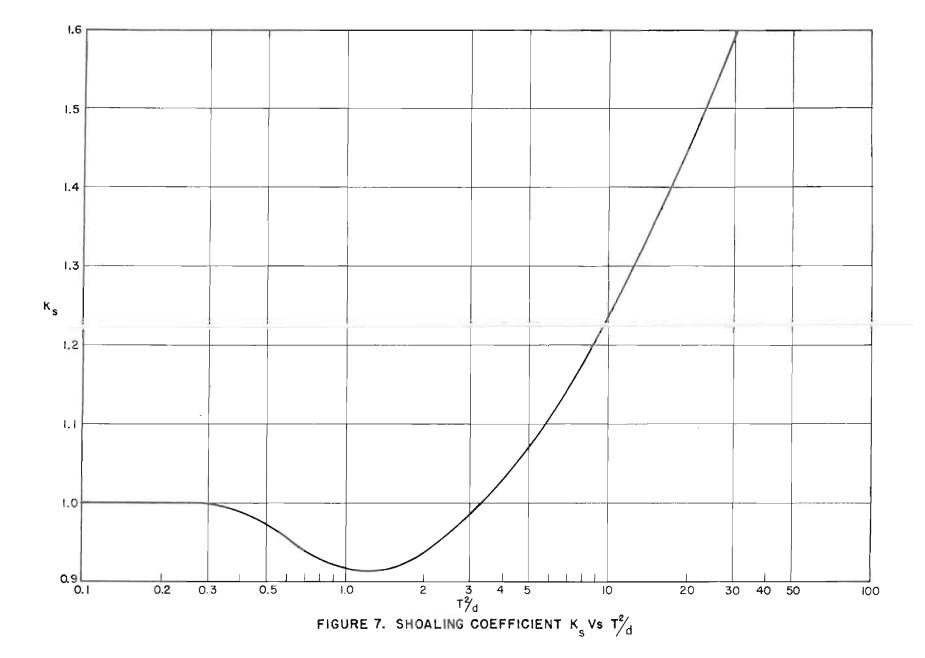
FOR COMPUTING WIND WAY ES IN SHALLOW WATER

- e. A deep water wave height, H and wave period T is determined at the beginning of the first increment of ΔF using Figure 5 and $F = F_{min}$.
- f. This value of H_0 is then assumed to travel over the increment ΔF as swell taking bottom friction into account. This is done by use of Figure 3. The quantity $\frac{fH\Delta X}{d^2}$ is determined using bottom friction factor f=.01 based on Lake Okeechobee studies and wave data in shallow water of the Gulf of Mexico; $H=K_{r}K_{0}$, where $K_{r}K_{0}$ is the shoaling factor (obtained from Figure 7); $K_{r}K_{0}$ is the refraction coefficient over the increment ΔF ; $\Delta X=6080^{\circ}\Delta F$, where ΔF is in nautical miles; $d=d_{r}K_{0}$, average depth over the increment ΔF ; and ΔF is computed using the average significant period over the increment and the average depth, $d_{r}K_{0}$ is read from Figure 3 and the actual significant height is equal to $H_{r}K_{0}$.
- g. An equivalent deep-water wave height H^{\bullet}_{o} is obtained from $H^{\bullet}_{o} = K_{f} H_{o}^{\bullet}$.
- h. Using H' and Figure 5 for example, an equivalent deepwater fetch length F' is obtained. For the case of regeneration of wind waves one also obtains an equivalent deep-water period, T_0 .
- i. An equivalent deep-water wave height is determined at the end of the second increment using $F = F^*_e + \Delta F = F_{\min}$, using Figure 5 for example. For the case of regeneration of wind waves one also obtains an equivalent deep-water period.
- j. With the average wave height 1/2 (H'_{ol} + H'_{ol}) steps f, g,h, and i are repeated. (This gives the swell height when the wave period T is held constant; T is given for U at F = F_{min}.) Using the average wave period 1/2 (T_{ol} + T_{ol}) steps f, g, h, and i are repeated. This gives the regenerated wind-vave height. The above procedure is used for all except the last increment or until the wave breaks, whichever occurs first. Usually the 1st increment cannot be treated by the above method since here the bettom slope increases too rapidly. The above procedure can be used for depths from deep water to about 20 feet, but has been applied to cepths of 12 feet, when the winds are not too great.

k. The maximum wave he: ght for all except the last increment is computed from

$$H_{\text{max}} = H_{\text{S}} \left[\frac{145 \text{ gd}}{\text{H}^2} \right]^{0.1} + 10\% \ge .78d \ge 1.87 \text{ H}_{\text{S}}$$
 (12)

The above formula is obtained from a report (13) based on wind-generated waves recorded in the Gulf of Mexico and hurricane wind-generated waves recorded in Lake Okeechobee, Florida.



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1. Maximum waves over the last increment and the location of breaking are obtained by plotting $H_b = .78d$ seaward from the coast. A continuation of the smooth curve of H_{max} obtained in (k) is made until it intersects the curve of $H_{\text{max}} = .78d$.

In general the swell heights (under the influence of wind) and the regenerated wind-wave heights in the intermediate to shallow water will be about the same, except in the reaking zone where the swell having a larger shoaling factor, becomes dominant, and the wind waves will be superimposed thereon. The reason swell heights (under the influence of wind) are about the same as the regenerated wind waves in intermediate to shallow water is that a issipation of wave energy by bottom friction tends to become independent of wave period. This can be seen from Figure 2 where the isolines recome crowded for large values of T2/d. It is for this reason that one can use the wave-period relationships given for deep water when computing generation of wind waves in shallow water.

As an example, the above procedure has been used to compute wind waves from deep water over the bottom profile perpendicular to shore off Caplen, Texas, using U = 26 kt ots and F_{min} = 200 nautical miles (over deep water to d/L = 0.5). The results of these computations are given in Table 1 and Figure 8. This particular profile was selected since it consists of several different locations having the same depth. It is quite an irregular bottom profile, others being less irregular. For example, from Table 1 and Figure 8 there are two 40-foot depths, one for which H' = 10.9 feet and the other H' = 8.7 feet, the average value of H' = 9.8 was used to corpute gH' /U20 versus gd/U2 and gF/U2 used in Figure 11. This shows that one can expect forecasts on the order of + 10% of the values obtained from the forecasting graphs, discussed below.

GENERALIZED FORECASTING CURVIS FOR THE GULF OF MEXICO

In general, for any location on the Continental Shelf, three special cases for wind-wave generation exist: (a) winds blowing parallel to the coast, (b) winds blowing from land to sea and (c) winds blowing from sea to land. A set of generalized forecasting curves has been prepared based on the above numerical procedures. These forecasting curves are discussed below:

- a. Case I, Winds Blowing Parallel to the Coast. In this case, except where very irregular bottom topography exists, the best approach is to use the flat bottom relationships of gH. /U2, versus gd/U2 and gF/U2 given in Figure 9.
- b. Case II, Winds Blowing from Land to Sea. In a report by Bretschneider and Thompson(14) it was shown that for most offshore winds, waves are generated which to not feel the bottom. This is

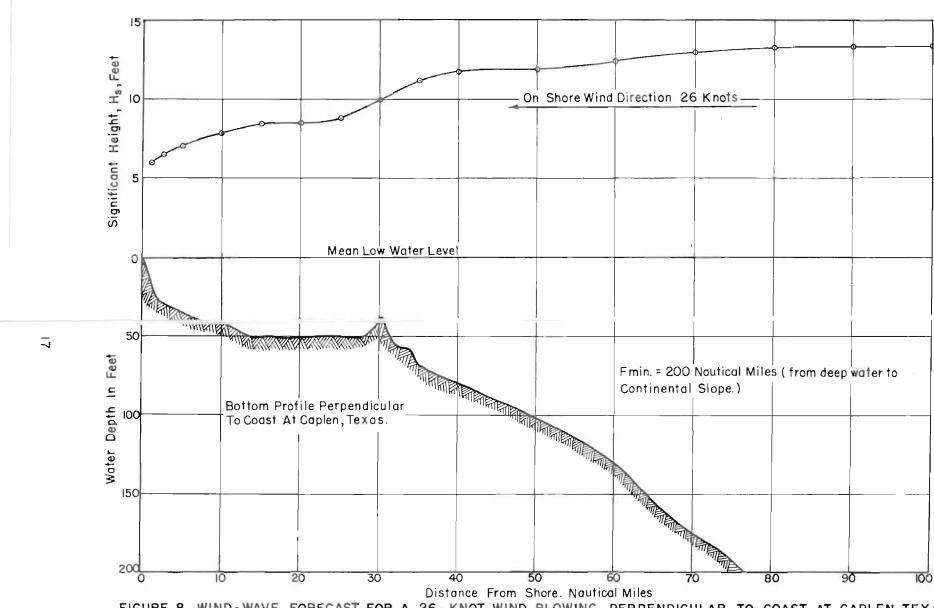


FIGURE 8. WIND-WAVE FORECAST FOR A 26- KNOT WIND BLOWING PERPENDICULAR TO COAST AT CAPLEN, TEXAS

TABLE I
SUMMARY OF COMPUTATIONS FOR EXAMPLE OF WIND-WAVE GENERATION OVER THE CONTINENTAL SHELF

	Line		Incre- ment_																	
	1	x*	n.m.	100	90	80	70	60	50	40	35	30	25	20	15	10	5	2.5	$\frac{16}{1.25}$	17
	2	P min	n.m.	200	200	200	190	187	182	155	136	112	74	53.5	48.5	45	39.5	29.5	25.2	
	3	۵F	n.m.	10	10	10	10	10	10	5	5	5	5	5	.5	5	5	2.5	1.25	1,25
	4	d ₁	feet	deep	300	230	190	175	130	100	80	67	40	50	51	51	38	34	28	20
	5	d ₂	feet	300	230	190	175	130	100	80	67	40	50	51	51	38	34	28	20	0
	ה	dave	feet	deep	265	210	182,5	152.5	115	90	73,5	51.5	45	50.5	51	44.5	36	31	24	
	7	H _{o1}	feet	13.3	13.3	13.3	13,29	13,28	13.21	12.94	12.66	12.19	10.89	9.67	9,35	9.09	8.65	7.72	7.30	
	8	H _{o2}	feet	13.3	13,3	13.3	13.30	13.30	13,29	13,00	12.76	12.31	11.11	10.01	9.67	9.43	9.03	8.00	7.42	
	9	Havé	feet	13.3	13,3	13.3	13,295	13,29	13.25	12.97	12.71	12.25	11.00	9.84	9.51	9.26	8.84	7.86	7.36	
	10	T_,	sec	9.4	9.4	9.4	9.3	9.3	9.24	8.94	8.50	8.25	7.52	6.94	6.76	6.6	6.38	5.82	5,52	
	11	T _{o2}	sec	9.4	9.4	9.4	9.4	9.4	9.40	9.00	8.76	8,35	7.62	7.10	6.94	6.8	6.60	6.00	5.62	
	12	To ave	sec	9.4	9.4	9.4	9,35	9.35	9,32	8.97	8.63	8.30	7.57	7.02	6.85	6.7	6.49	5.91	5.57	
8	13	$\frac{T^2}{d}$	sec ² /feet	deep	.333	.42	. 48	.572	.757	.894	1.01	1.335	1.29	.975	.920	1.01	1.17	1.12	1.29	
	14	K s ave		1.0	.996	.985	.975	,960	.932	.921	.918	.914	.913	.919	,920	,918	.913	.914	.913	
	15	fHAX d ²		Small	.114	.180	,212	.334	.567	. 438	. 656	1,285	1,51	1.09	1.02	1.30	1.89	1,14	.89	
	16	K _f		1.0	1.00	.9993	,9982	.9932	.974	.974	.955	.885	.870	.933	.940	.917	.855	.912	.92	
	7.2	H* o2	feet	13,3	13.3	13,29	13.28	13.21	12.94	12.66	12.19	10.89	9.67	9.35	9.09	8,65	7,72	7.30	6.83	
	18	Pe	n.m.	200	200	190	187	182	155	136	112	74	53.5	48.5	45	39.5	29.5	25.2	21.5	
	19	gH'o/U2	$(H^*_{0} = H^*_{02})$,221	.221	,221	.221	.219	.215	.210	.202	.181	.161	.155	.151	.144	.128	.121	.113	
	20	gd/U^2 (d	= d ₂)	4.98	3.82	3.16	2.91	2,16	1,66	1.33	1,11	.664	.831	.847	.847	.631	.565	.465	.322	i.

^{*} X is distance from shore.

because the fetch length, increasing leaward, is limiting. As the fetch length gets longer the wave per od gets longer, but the water depth becomes greater. This would in licate that for offshore winds, one may use the deep-water wave forecasting curves, Figure 2.

Actually Figure 9 can also be used, since the lines of constant gF/U² terminate on the right hand side of Figure 5 at deep water.

- c. Case III, Winds Blowing from Sea to Land. This is perhaps the most complex situation for we we generation, and no one set of curves can be developed similar to those for Cases I and II. Since each location has a different bottom rofile leading shoreward from various directions, a set of graphs would be required for each direction for each station. This becomes impraint ticable just as the numerical procedure discussed in the previous so ction becomes impracticable when one requires a multitude of forecasts forecasting graph for onshore winds in the Gulf of Mexico. The forecasting curves are computed by the procedure outlined above based on the following conditions:
- (1) The most frequent minimum fetch length was selected for each wind speed range and onshore direction, based on wind and fetch statistics using three years of A. H. Glenn and Associates (19). The most with a particular wind and direction vas determined from

$$F_{\min} = \frac{F_1 t_1 + F_2 t_2 + F_3 t_3 + \dots + F_n t_n}{t_1 + t_2 + t_3 + \dots t_n}$$
(13)

where F is the minimum fetch length and t the duration of a particular wind and direction over this fetch, etc.

- (2) For each location an average bottom profile was determined based on three shoreward by thom profiles, forming the cone of 45° on each side of the profile perpendicular to the coast. The profile perpendicular to the coast was given double weight in determining the average. This was done for each location, resulting in five average bottom profiles. Each of the average bottom profiles was divided into short increments, and an average depth over each increment was determined. Computation of wind waves were then made as outlined above.
 - (3) Values of H°_{o} , F_{\min} , d_{F} and U were tabulated, d_{F}

depth at end of fetch F where H' applies. From the above, values of gH' $_{0}$ /U², gF/U² and gd/U², were determined where F is same as F_{min} and d same as d_F. A generalized plot was made of gH' $_{0}$ /U² versus gd/U² and isolines of gF/U². This required a certain amount of smoothing to obtain the best overall picture. Scatter of the data, however, was not significant, being less than $\pm 10\%$ of the final selected curves. The

resulting families of curves for each location are given in Figures 10. 11, 12, 13 and 14 respectively for Brownsville, Texas; Caplen, Texas; Burrwood, Louisiana; Apalachicola, Florida; and Tampa, Florida. These figures can be used for forecasting wind waves for onshore winds perpendicular to the coast ±45°.

These forecasting curves are used to compute the shallow water wave statistics given in references (1, 2, 3, 4, 5) making use of the duration of individual wind speed ranges and the corresponding most frequent minimum fetches.

It must be emphasized that certain assumptions were required to obtain the forecasting curves for wave statistics purposes, and that individual forecasts may deviate by ±10%. It is therefore, necessary, that any individual forecast take these factors into consideration. Since the curves have not yet been used for daily forecasts, it is difficult to recommend any modification except that the following suggestions seem in order:

- (1) For onshore winds perpendicular to the coast one might multiply values of H* by 1.1, where H* is determined from the proper figure.
- (2) For onshore winds plowing at 45° to the coast one might multiply values of H° by 0.9. (The above rules of thumb need to be checked against observations.)
- (3) Allowance should be made for location ahead of and behind various shoal areas.

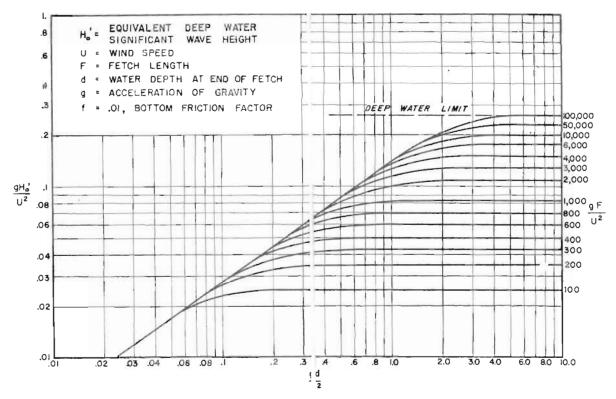
DETERMINATION OF SIGNI ICANT WAVE PERIOD

No generalized forecasting curves are presented for obtaining the significant wave period. Although a wave period is necessarily used and obtained in the computations for wave height it is not necessarily the significant wave period, although it may be close to the significant period. However, the significant wave period has a "definite significant relationship" to the significant height. It is shown by Bretschneider(20) for de p-water wind waves for gF/U² = 10,000 that the average value of the significant period is related to the significant height by:

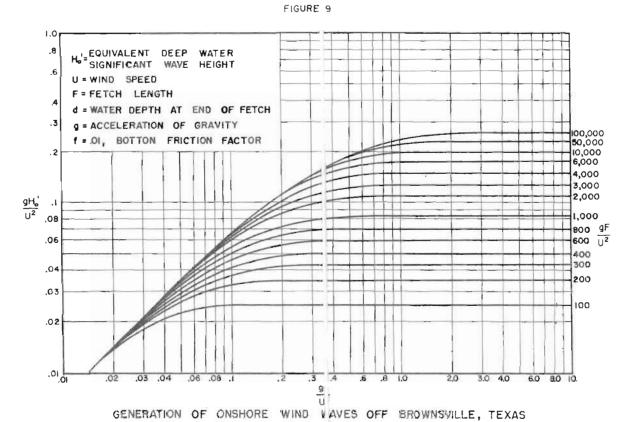
$$H_0/T_0^2 = 0.22$$
 (14)

This expression is quite signific and since it agrees with Jasper(21), wherein one may determine the average value of H /T 2 \cong 0.20, based on deep-water wave observations in the North Atlantic. From Figure 3 one may determine for shallow water (gd/U² = .01 to 1.0):

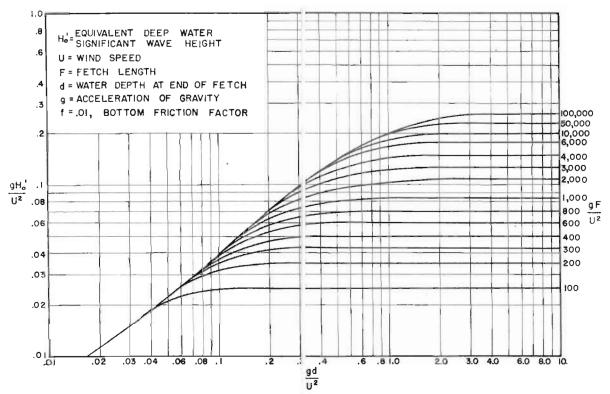
$$H^{\bullet}_{0}/T^{2} = .17 \pm 20\%$$
 (15)



GENERATION OF WIND WAVES OVER A BOTTOM OF CONSTANT DEP'IH FOR UNLIMITED WIND DURATION REPRESEN ED AS DIMENSIONLESS PARAMETERS

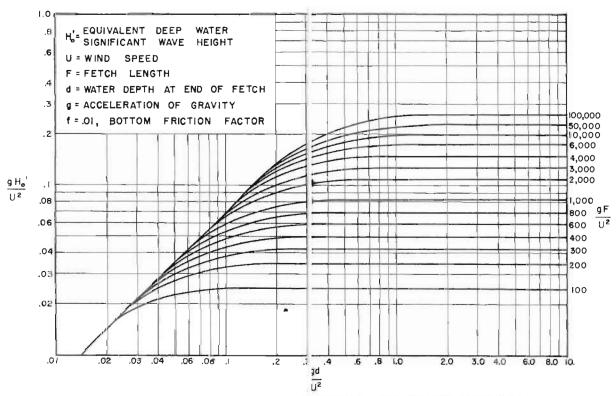


BASED ON AVERAGE CONDITIONS



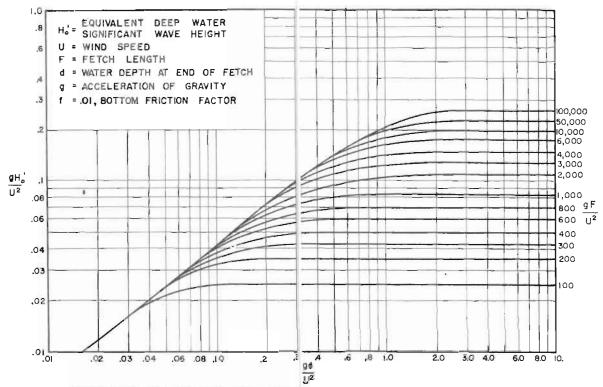
GENERATION OF ONSHORE WIND WAVES OFF CAPLEN, TEXAS BASED ON AVERAGE CONDITIONS

FIGL RE II



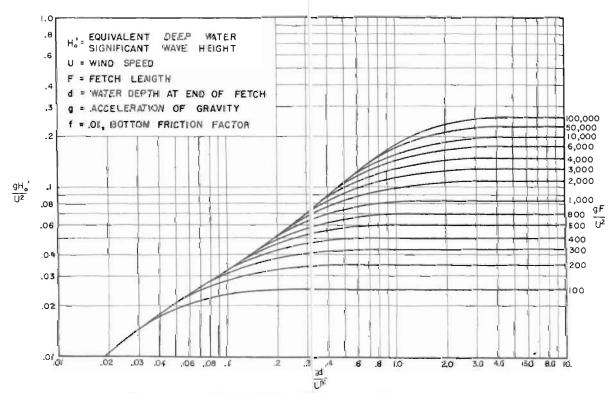
GENERATION OF ONSHORE WIND WAVES OFF BURRWOOD, LOUISIANA BASED ON AVE LAGE CONDITIONS

FIGU IE 12



GENERATION OF ONSHORE WIND WAYES OFF APALACHICOLA, FLORIDA BASED ON AVEFAGE CONDITIONS





GENERATION OF ONSHORE WIND WATES OFF TAMPA BAY, FLORIDA BASED ON AVERAGE CONDITIONS

FIGU ₹E 14

Using expressions 14 and 15 it follows:

$$T_0 = 2.13 \sqrt{H_0}$$
 for deep water (1b)

$$T_s = 2.43 \sqrt{H^*_o} + 10\%$$
 for shallow water (17)

Once the significant wave heights of wind waves have been obtained it is possible to obtain approximately the significant wave period directly from the above equations.

DISTRIBUTION FUNCTIONS FOR WIND WAYES FOR THE GULF OF MEXICO

The preceding material is used for height and period. Actual use of forecasts made from relationships presented in this report, as well as the in the Beach Brosion Board Technical Mem stake into account the distribution of all individual waves making up the wave train. Wind wave records to those of Putz(6), based on Pacific Ocean wave data. Therefore, Figures 15 and 16 reproduced from Putz(6) to obtain the distribution of heights and periods very close to those of Putz(6), based on Pacific Ocean wave data. Therefore, can be used temporarily to obtain the distribution of heights and periods, respectively, once the significant waves are forecast. Relationships between the significant wave height, H; average wave height, H apercent, H₁₀; and the maximum wave, H_{max}; average of highest ten the percent, M₁₀; and Siewell(17), based on the analysis of wave records. Snodgrass(18) summarizes the average values as follows:

$$H_{\text{max}}/H_{\text{s}} = 1.87 \pm 20\%$$
 (18)

$$H_{1/10}/H_s = 1.29 + 10\%$$
 (19)

$$H_{ave}/H_{s} = .65 + 10\%$$
 (20)

Based on the analysis of wind wave records in shallow water, Bretschneider(13) finds that:

$$H/H_{\text{max}} = \left[145 \frac{\text{gd}}{\text{U}^2}\right]^{-1} : 10\% \text{ for } \frac{\text{gd}}{\text{U}^2} = 2.5$$
 (21)

$$H_{1/10}/H_s = 1.23 \pm 10\%$$
 (22)

Additional analysis of these same data g ves:

$$H_{ave}/H_{s} = .675 \pm 10\%$$
 (23)

²⁴

^{*}Analysis of all wave records from the Gilf of Mexico has not yet been completed and the results are therefore not presented at this time.

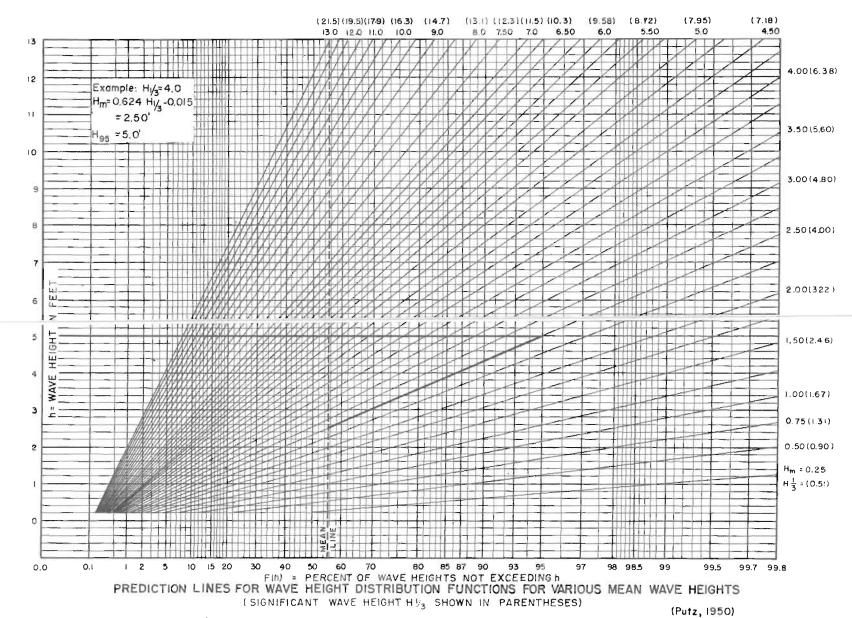
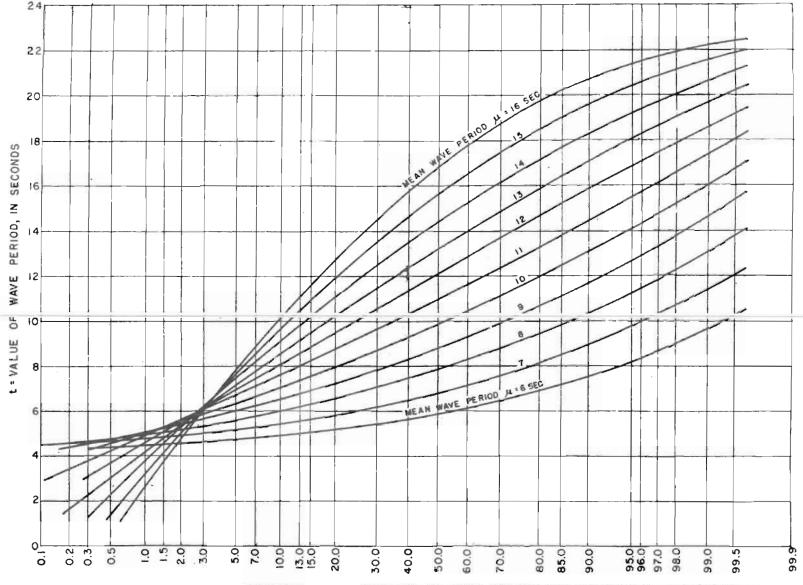


FIGURE 15



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APPENDIX \

EXAMPLE FORECIST

Determine the significant wave height, H_s; maximum probable wave height, H_{max}; average of highest ten percent of waves, H₁/10; and the average wave, H_{ave}; for the following conditions:

- (a) Location off Caplen, Texas
- (b) Minimum fetch length of 200 nautical miles (1,216,000 feet)
- (c) Unlimited wind durat on
- (d) Wind speed 26 knots 44 feet per/sec)
- (e) Wind direction, SE
- (I) for deep water
- (II) at a water depth of 0 feet
- (III) at a water depth of .: 0 feet
- (IV) at a water depth of 0 feet

Solution:

$$gF/U^2 = \frac{32.2 \times 1}{44} \frac{216,000}{2} = 20,300$$

I - for deep water

$$gd/U^2 = 6.6$$

$$\frac{gH'}{U^2} = ,22 \text{ (Fig. 7)}$$

$$H_0^2 = H_S = \frac{.22 (44)^2}{32.2} = 13.3 \text{ feet}$$

$$T_s = 2.13 \sqrt{13.3} (\pm 10\%) = 7.75 \text{ seconds } \pm 10\%$$

$$H_{\text{max}} = 13.3 \times 1.87 = 24.9 \text{ feet } \pm 20\% \text{ (equation 18)}$$

$$H_{1/10} = 13.3 \times 1.29 = 17.2 \text{ feet } \pm 10\% \text{ (equation 19)}$$

$$H_{ave} = 13.3 \text{ x .64} = 8.52 \text{ feet } \pm 10\% \text{ (equation 20)}$$

$$gd/U^2 = \frac{32.2 \times 40}{(44)^2} = .668$$

$$\frac{gH^*}{0} = .175 \text{ (Fig. 11)}$$

$$H^*_{0} = \frac{.175 \text{ (44)}^2}{32.2} = 10 \text{ 5 feet (Note } H^*_{0} = 10.9 \text{ and } 8.7 \text{ from Table 1)}$$

$$T_{g} = 2.43 \sqrt{10.5} + 10\% = 7.9 \text{ sec } \pm 10\% \text{ (equation } 17)$$

$$T_{g}^2/d = \frac{(7.9)^2}{40} = 1.5$$

$$K_{g} = .919 \text{ (Fig. 7)}$$

$$H_{g} = 10.8 \times .92 = 9.85 \text{ feet}$$

$$H_{max} = 9.85 \begin{bmatrix} 1.45 \times .66 \end{bmatrix} = 9.85 \times 1.58 = 15.5 \text{ feet}$$

$$\pm 10\% \text{ (equation } 21)$$

$$H_{1/10} = 9.85 \times 1.23 = 2.1 \text{ feet } \pm 10\% \text{ (equation } 22)$$

$$H_{ave} = 9.85 \times .675 = 6.65 \text{ feet } \pm 10\% \text{ (equation } 23)$$

$$III - \text{ for } d = 20 \text{ feet}$$

$$\frac{gd}{u^2} = \frac{32.2 \times 20}{(44)^2} = .33 \text{ (if } \frac{1}{32.2} = 6.5 \text{ feet (Note: } H^*_{0} = 6.83 \text{ Table 1)}$$

$$T_{g} = 2.43 \sqrt{6.5 \pm 10\%} = 6.10 \text{ seconds } \pm 10\% \text{ (equation } 17)$$

$$T_{g}^2/d = 1.92$$

$$K_{g} = .929 \text{ (Fig. 7)}$$

$$H_{g} = 6.3 \text{ x .} 929 = 5.85 \text{ feet}$$
 $H_{max} = 5.85 \begin{bmatrix} 145 \text{ x} \\ .334 \end{bmatrix}^{0.1} = 5.85 \text{ x } 1.47 = 8.6 \text{ feet}$
 $\pm 10\% \text{ (equation 21)}$
 $H_{1/10} = 5.85 \text{ x } 1.23 = 7.20 \text{ feet } \pm 10\% \text{ (equation 22)}$
 $H_{ave} = 5.85 \text{ x .} 675 = 3.96 \text{ feet } \pm 10\% \text{ (equation 23)}$
 $IV - \text{ for d} = 10 \text{ feet}^*$
 $\frac{gd}{u^2} = \frac{32.2 \text{ x } 10}{(44)^2} = 167$
 $H^*_{0} = .06 \frac{(44)^2}{32.2} = 3 = 6 \text{ feet}$
 $H_{0} = 0.06 \frac{(44)^2}{32.2} = 3 = 6 \text{ feet}$
 $H_{0} = 5.12 \begin{bmatrix} 4.15 \end{bmatrix}^2 = 108 \text{ feet}$
 $H_{0} = 10/108 = .0925$
 $H_{0} = .939 \text{ x } 3.6 = 3.4 \text{ feet}$
 $H_{max} = 3.4 \begin{bmatrix} 145 \text{ x .} 167 \end{bmatrix}^{-1} = 3.4 \text{ x } 1.375 = 4.8 \text{ feet}$
 $H_{1/10} = 3.4 \text{ x } 1.23 = 4.2 \text{ feet } \pm 10\% \text{ (equation 22)}$
 $H_{ave} = 3.4 \text{ x .} 675 = 2.3 \text{ feet } \pm 10\% \text{ (equation 23)}$

Δ-3

*Note: For this particular depth the inswers are only approximate, since the theory is extended too far into the steeper zone of the shoreward portion of the Continental Shelf. In fact, these waves may well be obscured by longer period swell breaking in the surf zone. The predominant period of the swell will be on the order of 7.9 seconds as given in the first example.